Mirror Symmetry of 3d $\mathcal{N} = 4$ Abelian Gauge Theories (joint with Andrew Ballin, Thomas Creutzig, Tudor Dimofte)

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Outline

1 Motivation

2 Boundary VOA

- A twist with Neumann b.c.
- B twist with Dirichlet b.c.
- Mirror symmetry of boundary VOA

3 The category of line operators

- The category of lines as modules of the boundary VOA
- Mirror symmetry of the category of line operators
- The effect of gauging

Future directions

Motivation

• A 3d $\mathcal{N} = 4$ abelian gauge theory: gauge group $(\mathbb{C}^*)^r$, hypermultiplets in \mathbb{C}^n , transform under charge matrix q • A 3d $\mathcal{N} = 4$ abelian gauge theory: gauge group $(\mathbb{C}^*)^r$, hypermultiplets in \mathbb{C}^n , transform under charge matrix $q \longrightarrow T_q$.

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- The category of line operators should be a braided tensor category (BTC).
- The work of Hilburn-Raskin gives chiral categories, but they are too big to possess braided tensor structure.

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- For $T_{A,q}$, \mathbb{B} is the Neumann boundary condition, and for $T_{B,q}$, \mathbb{B} is the Dirichlet boundary condition.
- Method was used in the previous work with A. Ballin.

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• This has level $-q^T q$:

$$J^{i}(z)J^{j}(w) \sim \frac{-\sum_{a} q_{ia}q_{ja}}{(z-w)^{2}}$$

5/21

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• We add $V_{bc}^{\otimes n}$ to cancel the anomaly, and the gauge group acts as:

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• The BRST differential

$$\tilde{Q}_{BRST} = \sum_{i} \oint \mathrm{d}z c^{i} \tilde{J}^{i}$$

Satisfies $\tilde{Q}^2 = 0$.

Definition

The boundary VOA $V_{A,q}$ for the Neumann boundary condition of $T_{A,q}$ is defined to be the cohomology:

$$V_{A,q} := H^* \left(V_{\beta\gamma}^{\otimes n} \otimes V_{bc}^{\otimes n} \otimes V_{bc}^{\otimes r}, \tilde{Q} \right)$$

• Costello-Gaiotto: perturbative VOA V_{per} . It has fields $N^i, E^i, \psi^{a,+}$ and $\psi^{a,-}$ with OPE:

$$N^{i}E^{j} \sim \frac{\delta_{ij}}{(z-w)^{2}}, \ N^{i}\psi^{a,\pm} = \frac{\pm q_{ai}\psi^{a,\pm}}{z-w},$$
$$\psi^{a,\pm}\psi^{b,-} \sim \frac{\delta_{ab}}{(z-w)^{2}} + \frac{\delta_{ab}\sum_{j}q_{aj}E^{j}}{z-w}.$$

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- Dirichlet boundary condition of gauge fields have boundary monopole operators. (Bullimore-Dimofte-Gaiotto-Hilburn) They correspond to modules of $V_{\rm per}$.

April 27, 2022

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 - The OPE:

$$N^{i}(z)|m\rangle \sim \frac{(q^{T}qm)_{i}/2}{z}, \ E^{i}(z)|m\rangle \sim \frac{m_{i}}{z}$$
$$\psi^{a,\pm}(z)|m\rangle \sim z^{\pm(qm)_{a}} : \psi^{a,\pm}(z)|m\rangle:$$
$$Y(|m\rangle, z)|n\rangle \sim z^{m^{T}q^{T}qn} : Y(|m\rangle, z)|n\rangle:$$

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$$\begin{split} N^{i}(z)|m\rangle &\sim \frac{(q^{T}qm)_{i}/2}{z}, \ E^{i}(z)|m\rangle \sim \frac{m_{i}}{z}\\ \psi^{a,\pm}(z)|m\rangle &\sim z^{\pm(qm)_{a}} :\psi^{a,\pm}(z)|m\rangle:\\ Y(|m\rangle,z)|n\rangle &\sim z^{m^{T}q^{T}qn} :Y(|m\rangle,z)|n\rangle: \end{split}$$

• We denote these simple modules by $V_{\text{per},m}$.

9/21

Definition (and claim)

We claim that $V_{per,m}$ are monopole operators, and define the boundary VOA for Dirichlet b.c. of $T_{B,q}$ to be:

$$V_{B,q} := \bigoplus_m V_{per,m}$$

Boundary VOA: mirror symmetry

$$0 \longrightarrow \mathbb{Z}^r \xrightarrow{q} \mathbb{Z}^n \xrightarrow{p} \mathbb{Z}^{n-r} \longrightarrow 0$$

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- $T_{A,q} \leftrightarrow T_{B,p^T}$.
- Expectation: Neumann on A corresponds to Dirichlet on B (Bullimore-Dimofte-Gaiotto-Hilburn).

Boundary VOA: mirror symmetry

Theorem (A. Ballin, T. Creutzig, T. Dimofte, W. N.)

There is an isomorphism of VOAs:

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Remark. This is obtained by using a free field realization of VOAs on both sides. For the left hand side, BRST cohomology of Fock modules of the Heisenberg VOA is well known.

The category of line operators: idea of construction

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- The problem: one needs to be careful when applying Huang-Lepowsky-Zhang to get braided tensor category (BTC).
- We can use the idea of simple current extensions (Creutzig-McRae-Yang, Creutzig-Kanade-Linshaw).

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- $KL_q^{[0]}$, the sub category of modules having trivial monodromy with $V_{B,q}$.

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 $M \times V_{\mathrm{per},m} \leftrightarrow M$

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Definition

We define the category of line operators for $T_{B,q}$, denoted by $\mathcal{C}_{B,q}$, to be $KL_q^{[0]}/\mathbb{Z}^r$.

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- Problems:
 - The category is not as explicit.
 - A gauging process is kept inexplicit.
- Another approach: $V_{A,q}$ is Morita equivalent to an extension of $V_{\beta\gamma}^{\otimes n}$.

• $V_{\beta\gamma}^{\otimes n}$ has \mathbb{Z}^n lattice of simple currents, given by:

$$\sigma^k V_{\beta\gamma}^{\otimes n} := \sigma^{k_1} V \otimes \sigma^{k_2} V \otimes \cdots \otimes \sigma^{k_n} V$$

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• Gauging chooses a sublattice such that k = qm. Extend:

$$V_{\mathrm{ext}} := \bigoplus_m \sigma^{qm} V_{\beta\gamma}^{\otimes n}.$$

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• We showed:

Proposition

There is an isomorphism of VOAs:

$$V_{ext} \otimes V_{bc}^{\otimes n} \cong V_{A,q} \otimes W$$

Here W and $V_{bc}^{\otimes n}$ have trivial category of modules.
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- Identify objects related by direct summands of V_{ext} :

 $M \leftrightarrow M \times \sigma^{qm} V_{\beta\gamma}^{\otimes n}.$

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Definition

We define the category of line operators $\mathcal{C}_{A,q}$ to be $\mathcal{C}_{\beta\gamma}^{\boxtimes n,[0],q}/\mathbb{Z}^r$.

17/21

The category of line operators: mirror symmetry

Theorem

There is an equivalence of braided tensor categories:

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Remark. The proof is based on the relation between $V_{\beta\gamma}$ and $V(\widehat{\mathfrak{gl}}(1|1))$ (w A. Ballin). All these come from subquotient of Kazhdan-Lusztig category of $V(\widehat{\mathfrak{gl}}(1|1))$.

The category of line operators: gauging

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- Recall that $\mathcal{K}^* = (1 + z\mathcal{O}) \times \mathbb{C}^* \times \mathbb{Z}$. Strongly equivariant D modules on \mathcal{K}^n .

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- Using $\mathcal{C}_{A,q}$, we can construct the coulomb branch for abelian gauge theories \longrightarrow what is the corresponding structure on the coulomb branch.
- What is the boundary VOA for the non-abelian gauge theories, and what is the category of line operators.

Thank you!